

## Deep Decomposition for Stochastic Normal-Abnormal Transport Supplementary Material

This supplementary material contains proofs for our representation theorems and additional implementation details for  $D^2$ -SONATA. Supp. **A** and Supp. **B** give the proofs for Theorem 1 and Theorem 2, respectively. Supp. **C** introduces our 2D/3D stochastic advection-diffusion PDE package in `PyTorch` and discusses numerical discretization and stability conditions. And Supp. **D** provides descriptions of how we generate our normal-abnormal brain advection-diffusion simulation dataset, including how we construct the velocity and diffusion fields with random abnormal patterns, and how we simulate the corresponding normal and abnormal brain advection-diffusion time-series.

### A. Theorem 1: Anomaly-decomposed Divergence-free Vector Representation

For any vector field  $\mathbf{V} \in L^p(\Omega)^d$  and scalar field  $A$  in  $\mathbb{R}_{(0,1]}(\Omega)$  on a bounded domain  $\Omega \subset \mathbb{R}^d$  with smooth boundary  $\partial\Omega$ . If  $\mathbf{V}$  satisfies  $\nabla \cdot \mathbf{V} = 0$ , there exist a potential  $\Psi$  in  $L^p(\Omega)^\alpha$  ( $\alpha = 1(3)$  when  $d = 2(3)$ ):

$$\mathbf{V} = \nabla \times (A\Psi), \quad (A\Psi) \cdot \mathbf{n}|_{\partial\Omega} = 0. \quad (12)$$

Conversely, for any  $A \in \mathbb{R}_{(0,1]}(\Omega)$ ,  $\Psi \in L^p(\Omega)^\alpha$ ,  $\nabla \cdot \mathbf{V} = \nabla \cdot (\nabla \times (A\Psi)) = 0$ .

(Here,  $L^p$  refers to the space of measurable functions for which the  $p$ -th power of the function absolute value is Lebesgue integrable. Specifically, let  $1 \leq p < \infty$  and  $(\Omega, \Sigma, \mu)$  be a measure space.  $L^p(\Omega)$  space is the set of all measurable functions whose absolute value raised to the  $p$ -th power has a finite integral, i.e.,  $\|f\|_p \equiv (\int_\Omega |f|^p d\mu)^{1/p} < \infty$ .)

*Proof.* The above Theorem is a Corollary of the following Theorem **A.1**, which is originally introduced by Liu et al. [37].

**Theorem A.1 (Divergence-free Vector Field Representation by the Curl of Potentials).** For any vector field  $\mathbf{V} \in L^p(\Omega)^d$  on a bounded domain  $\Omega \subset \mathbb{R}^d$  with smooth boundary  $\partial\Omega$ . If  $\mathbf{V}$  satisfies  $\nabla \cdot \mathbf{V} = 0$ , there exists a potential  $\Psi$  in  $L^p(\Omega)^\alpha$  such that ( $\alpha = 1$  when  $d = 2$ ,  $\alpha = 3$  when  $d = 3$ )

$$\mathbf{V} = \nabla \times \Psi, \quad \Psi \cdot \mathbf{n}|_{\partial\Omega} = 0, \quad \Psi \in L^p(\Omega)^\alpha. \quad (13)$$

Conversely, for any  $\Psi \in L^p(\Omega)^\alpha$ ,  $\nabla \cdot \mathbf{V} = \nabla \cdot (\nabla \times \Psi) = 0$ .

(For detailed proof of Theorem **A.1**, please refer to Supp. A in [37].)

Therefore, it is obvious that given any vector field  $\mathbf{V} \in L^p(\Omega)^d$  and scalar field  $A \in \mathbb{R}_{(0,1]}(\Omega)$  on a bounded domain  $\Omega \subset \mathbb{R}^d$  with smooth boundary  $\partial\Omega$ . If  $\mathbf{V}$  satisfies  $\nabla \cdot \mathbf{V} = 0$ , according to Theorem **A.1**, there exists a  $\tilde{\Psi}$  such that:

$$\begin{aligned} \mathbf{V} &\Leftrightarrow \nabla \times \tilde{\Psi} \\ &\Leftrightarrow \nabla \times \left( A \left( \tilde{\Psi}/A \right) \right) \\ &\Leftrightarrow \nabla \times (A\Psi) \end{aligned} \quad (14)$$

$$\begin{aligned} &\Leftrightarrow \nabla A \times \Psi + A \nabla \times \Psi \\ &\Leftrightarrow \nabla A \times \Psi + A \bar{\mathbf{V}}, \end{aligned} \quad (15)$$

where Eq. (14) corresponds to Eq. (12), and is the explicit expression for the implementation of ‘‘anomaly-encoded’’ velocity vector field  $\mathbf{V}$  in this paper.  $\bar{\mathbf{V}}$  in Eq. (15) refers to the ‘‘anomaly-free’’ velocity field as defined in Eq. (4) in the main manuscript where the relation between  $\mathbf{V}$  and  $\bar{\mathbf{V}}$  it is denoted as  $\mathbf{V} = A \diamond \bar{\mathbf{V}}$ . □

### B. Theorem 2: Anomaly-decomposed Symmetric PSD Tensor Representation

For any  $n \times n$  symmetric PSD tensor  $\mathbf{D}$  and  $A \in \mathbb{R}_{(0,1]}(\Omega)$ , there exist an upper triangular matrix with zero diagonal entries,  $\mathbf{B} \in \mathbb{R}^{\frac{n(n-1)}{2}}$ , and a non-negative diagonal matrix,  $\Lambda \in SD(n)$ , satisfying:

$$\mathbf{D} = \mathbf{U}(A\Lambda)\mathbf{U}^T, \quad \mathbf{U} = \exp(\mathbf{B} - \mathbf{B}^T) \in SO(n). \quad (16)$$

Conversely, for  $\forall A \in \mathbb{R}_{(0,1]}(\Omega)$ ,  $\forall \mathbf{B} \in \mathbb{R}^{\frac{n(n-1)}{2}}$ , and any diagonal matrix with non-negative diagonal entries,  $\Lambda \in SD(n)$ , Eq. (16) results in a symmetric PSD tensor,  $\mathbf{D}$ .

*Proof.* The above Theorem is a Corollary of the following Theorem B.1, which is originally introduced by Liu et al. [37].

**Theorem B.1 (Symmetric PSD Tensor Representation by Spectral Decomposition).** For any tensor  $\mathbf{D} \in PSD(n)$ , there exists an upper triangular matrix with zero diagonal entries,  $\mathbf{B} \in \mathbb{R}^{\frac{n(n-1)}{2}}$ , and a diagonal matrix with non-negative diagonal entries,  $\Lambda \in SD(n)$ , satisfying:

$$\mathbf{D} = \mathbf{U}\Lambda\mathbf{U}^T, \quad \mathbf{U} = \exp(\mathbf{B} - \mathbf{B}^T) \in SO(n). \quad (17)$$

Conversely, for any upper triangular matrix with zero diagonal entries,  $\mathbf{B} \in \mathbb{R}^{\frac{n(n-1)}{2}}$ , and any diagonal matrix with non-negative diagonal entries,  $\Lambda \in SD(n)$ , Eq. (17) results in a symmetric PSD tensor,  $\mathbf{D} \in PSD(n)$ .

(For detailed proof of Theorem B.1, please refer to Supp. B in [37].)

Therefore,  $\forall n \times n$  symmetric PSD tensors  $\mathbf{D}$  and  $\forall A \in \mathbb{R}_{(0,1]}(\Omega)$ , there exists a  $\mathbf{B} \in \mathbb{R}^{\frac{n(n-1)}{2}}$ , and a  $\tilde{\Lambda} \in SD(n)$ , such that:

$$\begin{aligned} \mathbf{D} &\Leftrightarrow \mathbf{U}\tilde{\Lambda}\mathbf{U}^T \\ &\Leftrightarrow \mathbf{U} \left( A \left( \tilde{\Lambda}/A \right) \right) \mathbf{U}^T \\ &\Leftrightarrow \mathbf{U}(A\Lambda)\mathbf{U}^T \\ &\Leftrightarrow A\mathbf{U}\Lambda\mathbf{U}^T \end{aligned} \quad (18)$$

$$\Leftrightarrow A\bar{\mathbf{D}}, \quad \mathbf{U} = \exp(\mathbf{B} - \mathbf{B}^T) \in SO(n), \quad (19)$$

where Eq. (18) corresponds to Eq. (16), and is the explicit expression for the implementation of ‘‘anomaly-encoded’’ diffusion tensor field  $\mathbf{D}$  in this paper.  $\bar{\mathbf{D}}$  in Eq. (19) refers to the ‘‘anomaly-free’’ diffusion tensor field as defined in Eq. (4) in the main manuscript where the relation between  $\mathbf{D}$  and  $\bar{\mathbf{D}}$  it is denoted as  $\mathbf{D} = A \circ \bar{\mathbf{D}}$ . □

## C. PyTorch Deterministic-Stochastic Advection-Diffusion PDE Toolkit: Numerical Derivations

Our stochastic advection-diffusion PDE toolkit is designed to solve, separately or jointly, deterministically or stochastically, advection and diffusion PDEs in 1D/2D/3D.

$$\begin{aligned} \frac{\partial C(\mathbf{x}, t)}{\partial t} &= \left. \frac{\partial C(\mathbf{x}, t)}{\partial t} \right|_{\text{adv.}} + \left. \frac{\partial C(\mathbf{x}, t)}{\partial t} \right|_{\text{diff.}} + \left. \frac{\partial C(\mathbf{x}, t)}{\partial t} \right|_{\text{sto.}} = \underbrace{-\nabla \cdot (\mathbf{V}(\mathbf{x}) \cdot C(\mathbf{x}, t))}_{\text{Fluid flow}} + \underbrace{\nabla \cdot (\mathbf{D}(\mathbf{x}) \nabla C(\mathbf{x}, t))}_{\text{Diffusion}} + \underbrace{\sigma(\mathbf{x}) \partial W(\mathbf{x}, t)}_{\text{Model Uncertainty}} \\ (\nabla \cdot \mathbf{V} = 0) &= \underbrace{-\mathbf{V}(\mathbf{x}) \cdot \nabla C(\mathbf{x}, t)}_{\text{Incompressible flow}} + \underbrace{\nabla \cdot (\mathbf{D}(\mathbf{x}) \nabla C(\mathbf{x}, t))}_{\text{Diffusion}} + \underbrace{\sigma(\mathbf{x}) \partial W(\mathbf{x}, t)}_{\text{Model Uncertainty}}. \end{aligned} \quad (20)$$

One can choose to model the velocity field as a constant, a general vector field, or a divergence-free vector field (for incompressible flow). Furthermore, the diffusion field can be modeled as a constant, a non-negative scalar field, or a symmetric positive semi-definite (PSD) tensor field. By controlling  $\sigma(\mathbf{x})$ , one can specify the variance the Brownian motion  $W(\mathbf{x}, t)$ , and thus the levels of uncertainty of the entire system. The toolkit is implemented as a custom `torch.nn.Module` subclass, such that one can directly use it as a (stochastic) advection-diffusion PDE solver for data simulation or easily wrap it into DNNs or numerical optimization frameworks for inverse (stochastic) PDE problems, i.e., parameters estimation.

### C.1. Advection and Diffusion Computation

The computation of advection and diffusion and their stability analysis are similar to YETI [37], please refer to Supp. C.1-C.2 in [37] for a detailed discussion.

## C.2. (Stochastic) Numerical Integration

As introduced above, after discretizing all the spatial derivatives on the right side of Eq. (20), we obtain a system of ordinary differential equations (ODEs), which can be solved by numerical integration [18].

If the system is modeled stochastically, the stochastic term,  $\sigma(\mathbf{x})\partial W(\mathbf{x}, t)$ , should be included during the integration.

**Definition 2 (Brownian motion).** A stochastic process  $(W_t)$  such that (1)  $W_0 = 0$ ; (2)  $(W_t - W_s) \sim \mathcal{N}(0, t - s)$ ,  $\forall t \geq s \geq 0$ ; (3) For all disjoint time interval pairs  $[t_1, t_2]$ ,  $[t_3, t_4]$  ( $t_1 < t_2 \leq t_3 \leq t_4$ ), the increments  $W_{t_4} - W_{t_3}$  and  $W_{t_2} - W_{t_1}$  are independent random variables.

We model the advection-diffusion process via a stochastic PDE (SPDE), where  $\sigma$  ( $\sigma(\mathbf{x}) \in \mathbb{R}$ ) denotes the variance of the Brownian motion  $W_t$  ( $W(\mathbf{x}, t) \in \mathbb{R}$ ) [4] and represents the epistemic uncertainty for the dynamical system. With this additional stochastic term, the existence and uniqueness of the solution to Eq. (20) still holds:

**Theorem 3 (Existence and uniqueness of SPDE [4, 49, 60]).** If the coefficients of the stochastic partial differential equation Eq. (20) with initial condition, satisfy the spatially-varying Lipschitz condition

$$\begin{aligned} & |\mathbf{V}(\mathbf{x}_1) - \mathbf{V}(\mathbf{x}_2)|^2 + |\mathbf{D}(\mathbf{x}_1) - \mathbf{D}(\mathbf{x}_2)|^2 \\ & + |\sigma(\mathbf{x}_1) - \sigma(\mathbf{x}_2)|^2 \leq K|\mathbf{x}_1 - \mathbf{x}_2|^2, \end{aligned} \quad (21)$$

and the spatial growth condition

$$|\mathbf{V}(\mathbf{x})|^2 + |\mathbf{D}(\mathbf{x})|^2 + |\sigma(\mathbf{x})|^2 \leq K(1 + \mathbf{x}^2), \quad (22)$$

then there is a continuous adapted solution  $C(\mathbf{x}, t)$  satisfying the  $L^2$  bound. Moreover, if  $C(\mathbf{x}, t)$  and  $\tilde{C}(\mathbf{x}, t)$  are both continuous solutions satisfying the  $L^2$  bound, then

$$P\left(C(\mathbf{x}, t) = \tilde{C}(\mathbf{x}, t) \text{ for all } t \in [0, T]\right) = 1. \quad (23)$$

For a detailed discussion regarding Theorem 3, please refer to [4, 49, 60]. During implementation, we use the Euler-Maruyama scheme [27, 28], and the discretized version of the stochastic term in Eq. (20) can therefore be written as  $\sigma(\mathbf{x})W(\mathbf{x}, t)/\sqrt{\Delta t}$ , where  $W \sim \mathcal{N}(0, 1)$ .

We then use the RK45 method to advance in time ( $\delta t$ ) to predict  $\hat{C}^{t+\delta t}$ . Note when the input mass transport time-series has relatively large temporal resolution ( $\Delta t$ ), the chosen  $\delta t$  should be smaller than  $\Delta t$  to satisfy the stability conditions (Supp. C.1), thereby ensuring stable numerical integration.

## D. Normal-Abnormal Brain Advection-Diffusion Dataset

Our brain advection-diffusion simulation dataset is based on the public IXI brain dataset<sup>8</sup>, from which we use 200 patients with complete collections of T1-/T2-weighted images, magnetic resonance angiography (MRA) images, and diffusion-weighted images (DWI) with 15 directions for the simulation of 3D divergence-free velocity vector and symmetric PSD diffusion tensor fields. All images above are resampled to isotropic spacing (1 mm), rigidly registered intra-subject (according to the MRA image), and brain-extracted using ITK<sup>9</sup>.

In general, the generation of the ‘‘anomaly-free’’ velocity vector fields and ‘‘anomaly-free’’ diffusion tensor fields are of the same procedures as in [37] (Please refer to Supp. D.1-D.2 in [37]). Here, Supp. D.1 provides the anomaly value fields (A) simulation proposed in this paper, and Supp. D.2 introduces the time-series simulation for normal-abnormal advection-diffusion processes.

### D.1. ‘‘Anomaly-encoded’’ Velocity Vector and Diffusion Tensor Fields (Fig. D.6)

For each case, we simulate both normal samples and anomaly-encoded samples. For samples treated as normal, we directly use the simulated  $\mathbf{V}$  and  $\mathbf{D}$  for the concentration time-series simulation. For anomaly encoding, the originally simulated  $\mathbf{V}$  and  $\mathbf{D}$  are treated as the ‘‘anomaly-free’’ fields for generating anomaly-encoded fields,  $\bar{\mathbf{V}}$  and  $\bar{\mathbf{D}}$ , via Eq. (4). For the generation of the anomaly value field  $A$ , its spatially varying value (within  $[0, 1]$ ) is determined by a Gaussian with center uniformly sampled across the entire spatial domain of the sample case.

<sup>8</sup>Available for download: <http://brain-development.org/ixi-dataset/>.

<sup>9</sup>Code in <https://github.com/InsightSoftwareConsortium/ITK>.

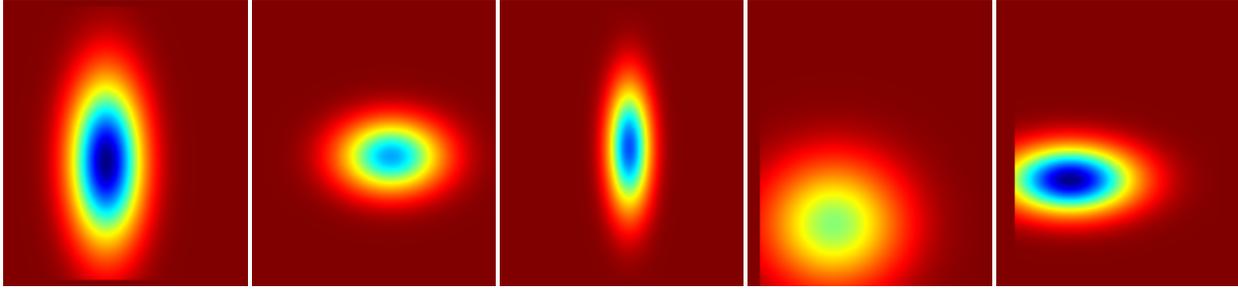


Figure D.6. Examples of different generated anomaly value field patterns applied on five different samples.

## D.2. Brain Advection-diffusion Time-series Simulation

For each brain advection-diffusion sample, the initial concentration state is assumed to be given by the MRA image with intensity ranges rescaled to  $[0, 1]$ . Time-series (length  $N_T = 40$ , interval  $\Delta t = 0.1 s$ ) are then simulated given the computed divergence-free velocity fields and symmetric PSD diffusion tensor fields by our advection-diffusion PDE solver (Supp. C). Thus the simulated dataset includes 800 brain advection-diffusion time-series (4 time-series for each of the 200 subjects, based on the four combinations of the simulated two velocity fields and two diffusion fields).