

# Supplementary Material

## DiffAvatar: Simulation-Ready Garment Optimization with Differentiable Simulation



Figure 1. After recovering simulation-ready assets, we can easily generate novel simulation results.

### 1. DiffXPBD: Differentiable Simulation

We provide details on the implementation of our differentiable simulator which builds upon DiffXPBD [1]. The simulation moves the states forward in time using  $\mathbf{q}_{n+1} = \mathbf{F}_n(\mathbf{q}_{n+1}, \mathbf{q}_n, \mathbf{u})$ , see Eq. (2). The adjoint states  $\hat{\mathbf{Q}}$  are computed in a backward pass using

$$\hat{\mathbf{q}}_{n-1} = \left( \frac{\partial \mathbf{F}_{n-1}}{\partial \mathbf{q}_n} \right)^\top \hat{\mathbf{q}}_{n-1} + \left( \frac{\partial \mathbf{F}_n}{\partial \mathbf{q}_n} \right)^\top \hat{\mathbf{q}}_n + \left( \frac{\partial \phi}{\partial \mathbf{q}_n} \right)^\top \quad (1)$$

The XPBD simulation frameworks uses the following update scheme.

$$\begin{aligned} \mathbf{x}_{n+1} &= \mathbf{x}_n + \Delta \mathbf{x}(\mathbf{x}_{n+1}) + \Delta t(\mathbf{v}_n + \Delta t \mathbf{M}^{-1} \mathbf{f}_{\text{ext}}) \\ \mathbf{v}_{n+1} &= \frac{1}{\Delta t}(\mathbf{x}_{n+1} - \mathbf{x}_n) \end{aligned} \quad (2)$$

We find the adjoint evolution for the XPBD integration scheme by combining this with (1) as

$$\begin{aligned} \hat{\mathbf{x}}_n &= \hat{\mathbf{x}}_{n+1} + \left( \frac{\partial \Delta \mathbf{x}}{\partial \mathbf{x}} + \Delta t^2 \mathbf{M}^{-1} \frac{\partial \mathbf{f}_{\text{ext}}}{\partial \mathbf{x}} \right)^\top \hat{\mathbf{x}}_n \\ &\quad + \frac{\hat{\mathbf{v}}_n}{\Delta t} - \frac{\hat{\mathbf{v}}_{n+1}}{\Delta t} + \frac{\partial \phi}{\partial \mathbf{x}}^\top \\ \hat{\mathbf{v}}_n &= \left( \frac{\partial \Delta \mathbf{x}}{\partial \mathbf{v}} + \Delta t^2 \mathbf{M}^{-1} \frac{\partial \mathbf{f}_{\text{ext}}}{\partial \mathbf{v}} \right)^\top \hat{\mathbf{x}}_n \\ &\quad + \Delta t \hat{\mathbf{x}}_{n+1} + \frac{\partial \phi}{\partial \mathbf{v}}^\top \end{aligned} \quad (3)$$

After re-arranging and by substituting  $\hat{\mathbf{v}}_n$  we find the adjoint states as

$$\begin{aligned} &\left( I - \frac{\partial \Delta \mathbf{x}}{\partial \mathbf{x}} - \Delta t^2 \mathbf{M}^{-1} \frac{\partial \mathbf{f}_{\text{ext}}}{\partial \mathbf{x}} - \frac{1}{\Delta t} \frac{\partial \Delta \mathbf{x}}{\partial \mathbf{v}} - \Delta t \mathbf{M}^{-1} \frac{\partial \mathbf{f}_{\text{ext}}}{\partial \mathbf{v}} \right)^\top \hat{\mathbf{x}}_n \\ &= 2 \hat{\mathbf{x}}_{n+1} - \frac{\hat{\mathbf{v}}_{n+1}}{\Delta t} + \frac{\partial \phi}{\partial \mathbf{x}}^\top + \frac{1}{\Delta t} \frac{\partial \phi}{\partial \mathbf{v}}^\top \end{aligned} \quad (4)$$

#### 1.1. Material Model

We use the orthotropic StVK model for modeling stretching and shearing and a hinge-based bending energy as detailed in [1]. The material parameters are recovered as part of the optimization process. Different material models can also be used.

### 2. Gradient of 3D Cloth Positions to 2D Patterns

To compute the gradient of the position with respect to the 2D patterns, we need to compute  $\frac{\partial \Delta \mathbf{x}}{\partial \bar{\mathbf{x}}_i}$ , for each of the 2D cloth vertex  $i \in [0, 1, \dots, n]$ . We use the same set of constraints as in DiffXPBD, where  $\mathbf{C} = [\epsilon_{00}, \epsilon_{11}, \epsilon_{01}]$ , and  $\epsilon$  is the Green strain.

$$\frac{\partial \Delta \mathbf{x}}{\partial \bar{\mathbf{x}}_i} = \mathbf{M}^{-1} \left( \frac{\partial \nabla \mathbf{C}}{\partial \bar{\mathbf{x}}_i} \Delta \lambda + \nabla \mathbf{C} \frac{\partial \Delta \lambda}{\partial \bar{\mathbf{x}}_i} \right) \quad (5)$$

$$\begin{aligned}
\frac{\partial \Delta \lambda}{\partial \bar{\mathbf{x}}_i} &= -\mathbf{J}^{-1} \frac{\partial \mathbf{J}}{\partial \bar{\mathbf{x}}_i} \mathbf{J}^{-1} \mathbf{b} + \mathbf{J}^{-1} \frac{\partial \mathbf{b}}{\partial \bar{\mathbf{x}}_i} \\
&= -\mathbf{J}^{-1} \left( \frac{\partial \mathbf{J}}{\partial \bar{\mathbf{x}}_i} \Delta \lambda - \frac{\partial \mathbf{b}}{\partial \bar{\mathbf{x}}_i} \right)
\end{aligned} \tag{6}$$

where  $\frac{\partial \mathbf{b}}{\partial \bar{\mathbf{x}}_i}$  and  $\frac{\partial \mathbf{J}}{\partial \bar{\mathbf{x}}_i} \Delta \lambda$  are computed as

$$\begin{aligned}
\frac{\partial \mathbf{b}}{\partial \bar{\mathbf{x}}_i} &= -\frac{\partial \mathbf{C}}{\partial \bar{\mathbf{x}}_i} - \frac{\partial \tilde{\alpha}}{\partial \bar{\mathbf{x}}_i} \lambda - \tilde{\alpha} \frac{\partial \lambda}{\partial \bar{\mathbf{x}}_i} \\
&= -\nabla \mathbf{C} - \tilde{\alpha} \sum \frac{\partial \Delta \lambda}{\partial \bar{\mathbf{x}}_i}
\end{aligned} \tag{7}$$

$$\begin{aligned}
\frac{\partial \mathbf{J}}{\partial \bar{\mathbf{x}}_i} \Delta \lambda &= \frac{\partial \nabla \mathbf{C}^T}{\partial \bar{\mathbf{x}}_i} \mathbf{M}^{-1} \nabla \mathbf{C} \Delta \lambda + \nabla \mathbf{C}^T \mathbf{M}^{-1} \frac{\partial \nabla \mathbf{C}}{\partial \bar{\mathbf{x}}_i} \Delta \lambda \\
&= \frac{\partial \nabla \mathbf{C}^T}{\partial \bar{\mathbf{x}}_i} \Delta \mathbf{x} + \nabla \mathbf{C}^T \mathbf{M}^{-1} \frac{\partial \nabla \mathbf{C}}{\partial \bar{\mathbf{x}}_i} \Delta \lambda
\end{aligned} \tag{8}$$

Given that  $\epsilon$  is a function of the deformation gradient  $\mathbf{F}$ , we provide the gradient of  $\mathbf{F}$  with respect to the rest positions, and the rest should just follow from chain rule. Note that  $\mathbf{F} = \mathbf{D} \bar{\mathbf{D}}^{-1}$ , where the columns of  $\mathbf{D}$  and  $\bar{\mathbf{D}}$  are the edge vectors, such that

$$\begin{aligned}
\mathbf{D} &= [\mathbf{x}_0 - \mathbf{x}_2 \quad \mathbf{x}_1 - \mathbf{x}_2] \\
\bar{\mathbf{D}} &= [\bar{\mathbf{x}}_0 - \bar{\mathbf{x}}_2 \quad \bar{\mathbf{x}}_1 - \bar{\mathbf{x}}_2]
\end{aligned} \tag{9}$$

The dimensions of the matrix  $\mathbf{D}$  is  $3 \times 2$ ,  $\bar{\mathbf{D}}$  is  $2 \times 2$ , and  $\mathbf{F}$  is  $3 \times 2$ .

We compute the derivative of the deformation gradient using the Einstein notation for  $\bar{\mathbf{x}}_0$  and  $\bar{\mathbf{x}}_1$

$$\begin{aligned}
\frac{\partial \mathbf{F}_{ij}}{\partial \bar{\mathbf{x}}_{mn}} &= \mathbf{D}_{ik} \frac{\partial \bar{\mathbf{D}}_{kj}^{-1}}{\partial \bar{\mathbf{x}}_{mn}} \\
&= -\mathbf{D}_{ik} \bar{\mathbf{D}}_{k\alpha}^{-1} \frac{\partial \bar{\mathbf{D}}_{\alpha\beta}^{-1}}{\partial \bar{\mathbf{x}}_{mn}} \bar{\mathbf{D}}_{\beta j}^{-1} \\
&= -\mathbf{D}_{ik} \bar{\mathbf{D}}_{km}^{-1} \bar{\mathbf{D}}_{nj}^{-1},
\end{aligned} \tag{10}$$

where  $\bar{\mathbf{x}}_{mn}$  is the  $n$ th component of  $\bar{\mathbf{x}}_m$ .

$$\frac{\partial \mathbf{F}_{ij}}{\partial \bar{\mathbf{x}}_2} = -\left( \frac{\partial \mathbf{F}_{ij}}{\partial \bar{\mathbf{x}}_0} + \frac{\partial \mathbf{F}_{ij}}{\partial \bar{\mathbf{x}}_1} \right) \tag{11}$$

### 3. Novel Animations

Figure 1 shows select frames from a novel simulated sequence with the recovered body shapes and garment patterns and materials.

### References

- [1] Tuur Stuyck and Hsiao-yu Chen. Diffxpbid: Differentiable position-based simulation of compliant constraint dynamics. *Proceedings of the ACM on Computer Graphics and Interactive Techniques*, 6(3):1–14, 2023. 1