

Appendices

A. Appendix A

We prove the invariance under local diffeomorphism that was used in Section 4 of the main paper and we then provide qualitative results on the NRSfM challenge dataset.

A.1. Invariance under Local Diffeomorphism

In figure 1, η is the image registration function, which means that we have $\mathbf{x} = \eta(\bar{\mathbf{x}})$. We can then write

$$\begin{pmatrix} du \\ dv \end{pmatrix} = \begin{pmatrix} \frac{\partial u}{\partial \bar{u}} & \frac{\partial u}{\partial \bar{v}} \\ \frac{\partial v}{\partial \bar{u}} & \frac{\partial v}{\partial \bar{v}} \end{pmatrix} \begin{pmatrix} d\bar{u} \\ d\bar{v} \end{pmatrix} = \mathbf{J}_\eta \begin{pmatrix} d\bar{u} \\ d\bar{v} \end{pmatrix}, \quad (23)$$

where \mathbf{J}_η is the Jacobian of the function η . Thus, under a change of variable from \mathbf{x} to $\bar{\mathbf{x}}$, we write $E(\phi \circ \eta)$ in terms of $E(\phi)$ (written as (2)) as

$$E(\phi \circ \eta) = ((\mathbf{e}_1 \quad \mathbf{e}_2) \mathbf{J}_\eta \quad |\mathbf{J}_\eta| \mathbf{e}_3) = E(\phi) \text{diag}(\mathbf{J}_\eta, |\mathbf{J}_\eta|). \quad (24)$$

We now introduce a theorem that describes the relation between moving frames related by a local diffeomorphism.

Theorem 1 (Moving frames under local diffeomorphism). *Given two surfaces \mathcal{S} and $\bar{\mathcal{S}}$ related by a local diffeomorphic mapping ψ , their respective moving frames E and \bar{E} are related by a linear transformation.*

Proof. From figure 1, we can write $\bar{\phi} = \psi \circ \phi \circ \eta$ and therefore,

$$\mathbf{J}_{\bar{\phi}} = \mathbf{J}_{\psi \circ \phi \circ \eta} \mathbf{J}_{\phi \circ \eta} \mathbf{J}_\eta. \quad (25)$$

According to inverse function theorem, \mathbf{J}_ψ is a linear function if ψ is locally diffeomorphic. Thus, we write $\mathbf{J}_\psi = \text{diag}(\lambda_1, \lambda_2, \lambda_3) \mathbf{R}$ where λ_i are scalars and \mathbf{R} is a rotation matrix. Using this result and the expression of $E(\phi)$ in terms of \mathbf{J}_ϕ in (2), we write (25) as

$$\bar{E}(\bar{\phi}) = \text{diag}(\lambda_1, \lambda_2, \lambda_3) \mathbf{R} E(\phi \circ \eta), \quad (26)$$

where $E(\phi \circ \eta)$ is given by (24). \square

Given the relation (25) obtained between moving frames under local diffeomorphism, we now derive the relation between connection components $\bar{\Gamma}_{jk}^i(\bar{\phi})$ and $\Gamma_{jk}^i(\phi)$ in the next theorem.

We write $w_j^i = \Gamma_{j1}^i du + \Gamma_{j2}^i dv$, and thus, the linear system in (3) can be written as

$$d\mathbf{e}_j = \frac{\partial \mathbf{e}_j}{\partial u} du + \frac{\partial \mathbf{e}_j}{\partial v} dv = w_j^1 \mathbf{e}_1 + w_j^2 \mathbf{e}_2 + w_j^3 \mathbf{e}_3, \quad j = [1, 2, 3] \quad (27)$$

Theorem 2 (Connection preservation under local diffeomorphism). *Given two surfaces \mathcal{S} and $\bar{\mathcal{S}}$ related by a local diffeomorphic mapping ψ , their respective connection components $\Gamma_{jk}^i(\phi)$ and $\bar{\Gamma}_{jk}^i(\bar{\phi})$ are preserved, i.e., $\bar{\Gamma}_{jk}^i(\bar{\phi}) = \Gamma_{jk}^i(\phi \circ \eta)$.*

Proof. Using the relation between \bar{E} and E obtained in (26) in (27), we get

$$\begin{pmatrix} \lambda_1 \bar{\mathbf{e}}_1^\top \\ \lambda_2 \bar{\mathbf{e}}_2^\top \\ \lambda_3 \bar{\mathbf{e}}_3^\top \end{pmatrix}^\top \mathbf{R}^\top \begin{pmatrix} \bar{w}_1^1 & \bar{w}_2^1 & \bar{w}_3^1 \\ \bar{w}_1^2 & \bar{w}_2^2 & \bar{w}_3^2 \\ \bar{w}_1^3 & \bar{w}_2^3 & \bar{w}_3^3 \end{pmatrix} = \begin{pmatrix} \lambda_1 \mathbf{e}_1^\top \\ \lambda_2 \mathbf{e}_2^\top \\ \lambda_3 \mathbf{e}_3^\top \end{pmatrix}^\top \mathbf{R}^\top \text{diag}(d\mathbf{J}_\eta, d|\mathbf{J}_\eta|) + \begin{pmatrix} \lambda_1 \mathbf{e}_1^\top \\ \lambda_2 \mathbf{e}_2^\top \\ \lambda_3 \mathbf{e}_3^\top \end{pmatrix}^\top \mathbf{R}^\top \begin{pmatrix} w_1^1 & w_2^1 & w_3^1 \\ w_1^2 & w_2^2 & w_3^2 \\ w_1^3 & w_2^3 & w_3^3 \end{pmatrix} \text{diag}(\mathbf{J}_\eta, |\mathbf{J}_\eta|). \quad (28)$$

The above relation is thus independent of λ_i and \mathbf{R} . We write $\bar{w}_j^i(\bar{\phi}) = w_j^i(\phi \circ \eta)$ and thus, $\bar{\Gamma}_{jk}^i(\bar{\phi}) = \Gamma_{jk}^i(\phi \circ \eta)$. \square

Upon expanding the relation between the connections (28), we get

$$\begin{pmatrix} \bar{\Gamma}_{11}^1 & \bar{\Gamma}_{11}^2 & \bar{\Gamma}_{11}^3 \\ \bar{\Gamma}_{21}^1 & \bar{\Gamma}_{21}^2 & \bar{\Gamma}_{21}^3 \\ \bar{\Gamma}_{31}^1 & \bar{\Gamma}_{31}^2 & \bar{\Gamma}_{31}^3 \end{pmatrix} = \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}^{-1} \begin{pmatrix} \frac{\partial \mathbf{J}_\eta}{\partial \bar{u}} & 0 \\ 0 & \frac{\partial |\mathbf{J}_\eta|}{\partial \bar{u}} \end{pmatrix} + \frac{\partial u}{\partial \bar{u}} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}^{-1} \begin{pmatrix} \Gamma_{11}^1 & \Gamma_{11}^2 & \Gamma_{11}^3 \\ \Gamma_{21}^1 & \Gamma_{21}^2 & \Gamma_{21}^3 \\ \Gamma_{31}^1 & \Gamma_{31}^2 & \Gamma_{31}^3 \end{pmatrix} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix} + \frac{\partial v}{\partial \bar{u}} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}^{-1} \begin{pmatrix} \Gamma_{12}^1 & \Gamma_{12}^2 & \Gamma_{12}^3 \\ \Gamma_{22}^1 & \Gamma_{22}^2 & \Gamma_{22}^3 \\ \Gamma_{32}^1 & \Gamma_{32}^2 & \Gamma_{32}^3 \end{pmatrix} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}, \begin{pmatrix} \bar{\Gamma}_{12}^1 & \bar{\Gamma}_{12}^2 & \bar{\Gamma}_{12}^3 \\ \bar{\Gamma}_{22}^1 & \bar{\Gamma}_{22}^2 & \bar{\Gamma}_{22}^3 \\ \bar{\Gamma}_{32}^1 & \bar{\Gamma}_{32}^2 & \bar{\Gamma}_{32}^3 \end{pmatrix} = \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}^{-1} \begin{pmatrix} \frac{\partial \mathbf{J}_\eta}{\partial \bar{v}} & 0 \\ 0 & \frac{\partial |\mathbf{J}_\eta|}{\partial \bar{v}} \end{pmatrix} + \frac{\partial u}{\partial \bar{v}} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}^{-1} \begin{pmatrix} \Gamma_{11}^1 & \Gamma_{11}^2 & \Gamma_{11}^3 \\ \Gamma_{21}^1 & \Gamma_{21}^2 & \Gamma_{21}^3 \\ \Gamma_{31}^1 & \Gamma_{31}^2 & \Gamma_{31}^3 \end{pmatrix} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix} + \frac{\partial v}{\partial \bar{v}} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}^{-1} \begin{pmatrix} \Gamma_{12}^1 & \Gamma_{12}^2 & \Gamma_{12}^3 \\ \Gamma_{22}^1 & \Gamma_{22}^2 & \Gamma_{22}^3 \\ \Gamma_{32}^1 & \Gamma_{32}^2 & \Gamma_{32}^3 \end{pmatrix} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}. \quad (29)$$

We write $\mathbf{J}_\eta = \begin{pmatrix} a & c \\ b & d \end{pmatrix}$ and the above expression can be written as

$$\begin{pmatrix} \bar{\Gamma}_{11}^1 & \bar{\Gamma}_{11}^2 & \bar{\Gamma}_{11}^3 \\ \bar{\Gamma}_{21}^1 & \bar{\Gamma}_{21}^2 & \bar{\Gamma}_{21}^3 \\ \bar{\Gamma}_{31}^1 & \bar{\Gamma}_{31}^2 & \bar{\Gamma}_{31}^3 \end{pmatrix} = \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}^{-1} \left(\begin{pmatrix} \frac{\partial \mathbf{J}_\eta}{\partial \bar{u}} & 0 \\ 0 & \frac{\partial |\mathbf{J}_\eta|}{\partial \bar{u}} \end{pmatrix} + \begin{pmatrix} a\Gamma_{11}^1 + b\Gamma_{12}^1 & a\Gamma_{11}^2 + b\Gamma_{12}^2 & a\Gamma_{11}^3 + b\Gamma_{12}^3 \\ a\Gamma_{21}^1 + b\Gamma_{22}^1 & a\Gamma_{21}^2 + b\Gamma_{22}^2 & a\Gamma_{21}^3 + b\Gamma_{22}^3 \\ a\Gamma_{31}^1 + b\Gamma_{32}^1 & a\Gamma_{31}^2 + b\Gamma_{32}^2 & a\Gamma_{31}^3 + b\Gamma_{32}^3 \end{pmatrix} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix} \right),$$

$$\begin{pmatrix} \bar{\Gamma}_{12}^1 & \bar{\Gamma}_{12}^2 & \bar{\Gamma}_{12}^3 \\ \bar{\Gamma}_{22}^1 & \bar{\Gamma}_{22}^2 & \bar{\Gamma}_{22}^3 \\ \bar{\Gamma}_{32}^1 & \bar{\Gamma}_{32}^2 & \bar{\Gamma}_{32}^3 \end{pmatrix} = \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix}^{-1} \left(\begin{pmatrix} \frac{\partial \mathbf{J}_\eta}{\partial \bar{v}} & 0 \\ 0 & \frac{\partial |\mathbf{J}_\eta|}{\partial \bar{v}} \end{pmatrix} + \begin{pmatrix} c\Gamma_{11}^1 + d\Gamma_{12}^1 & c\Gamma_{11}^2 + d\Gamma_{12}^2 & c\Gamma_{11}^3 + d\Gamma_{12}^3 \\ c\Gamma_{21}^1 + d\Gamma_{22}^1 & c\Gamma_{21}^2 + d\Gamma_{22}^2 & c\Gamma_{21}^3 + d\Gamma_{22}^3 \\ c\Gamma_{31}^1 + d\Gamma_{32}^1 & c\Gamma_{31}^2 + d\Gamma_{32}^2 & c\Gamma_{31}^3 + d\Gamma_{32}^3 \end{pmatrix} \begin{pmatrix} \mathbf{J}_\eta & 0 \\ 0 & |\mathbf{J}_\eta| \end{pmatrix} \right).$$

Substituting Γ_{jk}^i and $\bar{\Gamma}_{jk}^i$ from (4) gives

$$\begin{aligned} \bar{\Gamma}_{11}^1 &= \frac{2\bar{x}_1\bar{\beta}^4}{D} (1 + (1 + \bar{v}^2)\bar{x}_2^2 + \bar{u}\bar{x}_1 + 2\bar{v}\bar{x}_2 + \bar{u}\bar{v}\bar{x}_1\bar{x}_2) = \frac{2\beta^4(ax_1 + bx_2)d}{D\mathbf{J}_\eta} (a(1 + vx_2)^2 + ax_2^2 - b(1 + u^2)x_1x_2 + u(1 + vx_2)(ax_1 - bx_2)) - \\ &\frac{2\beta^4(ax_1 + bx_2)c}{D\mathbf{J}_\eta} (b(1 + ux_1)^2 + bx_1^2 - a(1 + v^2)x_1x_2 - v(1 + ux_1)(ax_1 - bx_2)) + \frac{1}{\mathbf{J}_\eta} \left(d\frac{\partial a}{\partial \bar{u}} - c\frac{\partial b}{\partial \bar{u}} \right), \\ \bar{\Gamma}_{11}^2 &= \frac{2\bar{x}_1^2\bar{\beta}^4}{D} (\bar{v} + (1 + \bar{v}^2)\bar{x}_2^2 + \bar{u}\bar{v}\bar{x}_1) = \frac{2\beta^4(ax_1 + bx_2)b}{D\mathbf{J}_\eta} (a(1 + vx_2)^2 + ax_2^2 - b(1 + u^2)x_1x_2 + u(1 + vx_2)(ax_1 - bx_2)) - \\ &\frac{2\beta^4(ax_1 + bx_2)d}{D\mathbf{J}_\eta} (b(1 + ux_1)^2 + bx_1^2 - a(1 + v^2)x_1x_2 - v(1 + ux_1)(ax_1 - bx_2)) + \frac{1}{\mathbf{J}_\eta} \left(b\frac{\partial a}{\partial \bar{u}} - a\frac{\partial b}{\partial \bar{u}} \right), \\ \bar{\Gamma}_{12}^1 &= \bar{\Gamma}_{21}^1 = \frac{\bar{x}_2\bar{\beta}^4}{D} (1 + (1 + \bar{v}^2)\bar{x}_2^2 - (1 + \bar{u}^2)\bar{x}_1^2 + 2\bar{v}\bar{x}_2) = \frac{\beta^4(ax_1 + bx_2)d}{D\mathbf{J}_\eta} (c(1 + vx_2)^2 + cx_2^2 - d(1 + u^2)x_1x_2 + u(1 + vx_2)(cx_1 - dx_2)) - \\ &\frac{\beta^4(ax_1 + bx_2)c}{D\mathbf{J}_\eta} (d(1 + ux_1)^2 + dx_1^2 - c(1 + v^2)x_1x_2 - u(1 + vx_2)(ax_1 - bx_2)) + \\ &\frac{\beta^4(cx_1 + dx_2)d}{D\mathbf{J}_\eta} (a(1 + vx_2)^2 + ax_2^2 - b(1 + u^2)x_1x_2 + u(1 + vx_2)(ax_1 - bx_2)) - \\ &\frac{\beta^4(cx_1 + dx_2)c}{D\mathbf{J}_\eta} (b(1 + ux_1)^2 + bx_1^2 - a(1 + v^2)x_1x_2 - v(1 + ux_1)(cx_1 - dx_2)) + \frac{1}{\mathbf{J}_\eta} \left(d\frac{\partial c}{\partial \bar{u}} - c\frac{\partial d}{\partial \bar{u}} \right), \\ \bar{\Gamma}_{12}^2 &= \bar{\Gamma}_{21}^2 = \frac{\bar{x}_1\bar{\beta}^4}{D} (1 + (1 + \bar{u}^2)\bar{x}_1^2 - (1 + \bar{v}^2)\bar{x}_2^2 + 2\bar{u}\bar{x}_1) = \frac{\beta^4(ax_1 + bx_2)a}{D\mathbf{J}_\eta} (d(1 + ux_1)^2 + dx_1^2 - c(1 + v^2)x_1x_2 - v(1 + ux_1)(ax_1 - bx_2)) - \\ &\frac{\beta^4(ax_1 + bx_2)a}{D\mathbf{J}_\eta} (c(1 + vx_2)^2 + cx_2^2 - d(1 + u^2)x_1x_2 + u(1 + vx_2)(cx_1 - dx_2)) + \\ &\frac{\beta^4(cx_1 + dx_2)a}{D\mathbf{J}_\eta} (b(1 + ux_1)^2 + bx_1^2 - a(1 + v^2)x_1x_2 - v(1 + ux_1)(cx_1 - dx_2)) - \\ &\frac{\beta^4(cx_1 + dx_2)b}{D\mathbf{J}_\eta} (a(1 + vx_2)^2 + cx_2^2 - d(1 + u^2)x_1x_2 + u(1 + vx_2)(ax_1 - bx_2)) - \frac{1}{\mathbf{J}_\eta} \left(b\frac{\partial c}{\partial \bar{u}} - a\frac{\partial d}{\partial \bar{u}} \right), \\ \bar{\Gamma}_{22}^1 &= -\frac{2\bar{x}_2^2\bar{\beta}^4}{D} (\bar{u} + (1 + \bar{u}^2)\bar{x}_1^2 + \bar{u}\bar{v}\bar{x}_2) = \frac{2\beta^4(cx_1 + dx_2)d}{D\mathbf{J}_\eta} (c(1 + vx_2)^2 + cx_2^2 - d(1 + u^2)x_1x_2 + u(1 + vx_2)(cx_1 - dx_2)) - \\ &\frac{2\beta^4(cx_1 + dx_2)c}{D\mathbf{J}_\eta} (d(1 + ux_1)^2 + dx_1^2 - c(1 + v^2)x_1x_2 - v(1 + ux_1)(cx_1 - dx_2)) + \frac{1}{\mathbf{J}_\eta} \left(d\frac{\partial c}{\partial \bar{v}} - c\frac{\partial d}{\partial \bar{v}} \right), \\ \bar{\Gamma}_{22}^2 &= \frac{2\bar{x}_2\bar{\beta}^4}{D} (1 + (1 + \bar{u}^2)\bar{x}_1^2 + \bar{v}\bar{x}_2 + 2\bar{u}\bar{x}_1 + \bar{u}\bar{v}\bar{x}_1\bar{x}_2) = \frac{2\beta^4(cx_1 + dx_2)a}{D\mathbf{J}_\eta} (d(1 + ux_1)^2 + dx_1^2 - c(1 + v^2)x_1x_2 - v(1 + ux_1)(cx_1 - dx_2)) - \\ &\frac{2\beta^4(cx_1 + dx_2)b}{D\mathbf{J}_\eta} (c(1 + vx_2)^2 + cx_2^2 - d(1 + u^2)x_1x_2 + u(1 + vx_2)(cx_1 - dx_2)) + \frac{1}{\mathbf{J}_\eta} \left(a\frac{\partial d}{\partial \bar{v}} - b\frac{\partial c}{\partial \bar{v}} \right), \\ \bar{\Gamma}_{31}^3 &= \bar{x}_1 + \frac{2\bar{\beta}^4\bar{x}_1}{D} (1 + \bar{u}\bar{x}_1 + \bar{v}\bar{x}_2) = (ax_1 + bx_2) \left(1 + \frac{2\beta^4}{D} (1 + ux_1 + vx_2) \right) + |\mathbf{J}_\eta|^{-1} \frac{\partial |\mathbf{J}_\eta|}{\partial \bar{u}}, \\ \bar{\Gamma}_{32}^3 &= \bar{x}_2 + \frac{2\bar{\beta}^4\bar{x}_2}{D} (1 + \bar{u}\bar{x}_1 + \bar{v}\bar{x}_2) = (cx_1 + dx_2) \left(1 + \frac{2\beta^4}{D} (1 + ux_1 + vx_2) \right) + |\mathbf{J}_\eta|^{-1} \frac{\partial |\mathbf{J}_\eta|}{\partial \bar{v}}, \\ \bar{\Gamma}_{11}^3 &= \frac{\bar{x}_1^2\bar{\beta}^3}{D} = \frac{\beta^3(ax_1 + bx_2)^2}{|\mathbf{J}_\eta|D}, \end{aligned}$$

$$\begin{aligned}
\bar{\Gamma}_{12}^3 &= \bar{\Gamma}_{21}^3 = \frac{\bar{x}_1 \bar{x}_2 \bar{\beta}^3}{\bar{D}} = \frac{\beta^3(ax_1+bx_2)(cx_1+dx_2)}{|\mathbf{J}_\eta|D}, \\
\bar{\Gamma}_{22}^3 &= \frac{\bar{x}_2^2 \bar{\beta}^3}{\bar{D}} = \frac{\beta^3(cx_1+dx_2)^2}{|\mathbf{J}_\eta|D}, \\
\bar{\Gamma}_{31}^1 &= \frac{\bar{\beta}^5(\bar{x}_1^2 + \bar{v}\bar{x}_1^2\bar{x}_2 - \bar{u}\bar{x}_1\bar{x}_2^2)}{\bar{D}} = \frac{\beta^5(ax_1+bx_2)}{D} (d(x_1+vx_1x_2-ux_2^2) - c(x_2+ux_1x_2-vx_1^2)), \\
\bar{\Gamma}_{31}^2 &= \frac{\bar{\beta}^5(\bar{x}_1\bar{x}_2 + \bar{u}\bar{x}_1^2\bar{x}_2 - \bar{v}\bar{x}_1^3)}{\bar{D}} = \frac{\beta^5(ax_1+cx_2)}{D} (-b(x_1+vx_1x_2-ux_2^2) + a(x_2+ux_1x_2-vx_1^2)), \\
\bar{\Gamma}_{32}^1 &= \frac{\bar{\beta}^5(\bar{x}_1\bar{x}_2 + \bar{v}\bar{x}_1\bar{x}_2^2 - \bar{u}\bar{x}_2^3)}{\bar{D}} = \frac{\beta^5(cx_1+dx_2)}{D} (d(x_1+vx_1x_2-ux_2^2) - c(x_2+ux_1x_2-vx_1^2)), \\
\bar{\Gamma}_{32}^2 &= \frac{\bar{\beta}^5(\bar{x}_2^2 + \bar{u}\bar{x}_1\bar{x}_2^2 - \bar{v}\bar{x}_1^2\bar{x}_2)}{\bar{D}} = \frac{\beta^5(cx_1+dx_2)}{D} (-b(x_1+vx_1x_2-ux_2^2) + a(x_2+ux_1x_2-vx_1^2)).
\end{aligned}$$

From the last expressions, we can that the last 8 connections $(\bar{\Gamma}_{11}^3, \bar{\Gamma}_{12}^3, \bar{\Gamma}_{12}^3, \bar{\Gamma}_{21}^3, \bar{\Gamma}_{22}^3, \bar{\Gamma}_{31}^1, \bar{\Gamma}_{31}^2, \bar{\Gamma}_{32}^1, \bar{\Gamma}_{32}^2)$, do not contain second order derivatives of η .